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Updated measurement of decay-time-dependent CP asymmetries in $D^0 \rightarrow K^+ K^-$ and $D^0 \rightarrow \pi^+ \pi^-$ decays

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A search for decay-time-dependent charge-parity (CP) asymmetry in $D^0 \rightarrow K^+ K^-$ and $D^0 \rightarrow \pi^+ \pi^-$ decays is performed at the LHCb experiment using proton-proton collision data recorded at a center-of-mass energy of 13 TeV, and corresponding to an integrated luminosity of 5.4 fb^{-1} . The D^0 mesons are required to originate from semileptonic decays of b hadrons, such that the charge of the muon identifies the flavor of the neutral D meson at production. The asymmetries in the effective decay widths of D^0 and \bar{D}^0 mesons are determined to be $A_\Gamma(K^+ K^-) = (-4.3 \pm 3.6 \pm 0.5) \times 10^{-4}$ and $A_\Gamma(\pi^+ \pi^-) = (2.2 \pm 7.0 \pm 0.8) \times 10^{-4}$, where the uncertainties are statistical and systematic, respectively. The results are consistent with CP symmetry and, when combined with previous LHCb results, yield $A_\Gamma(K^+ K^-) = (-4.4 \pm 2.3 \pm 0.6) \times 10^{-4}$ and $A_\Gamma(\pi^+ \pi^-) = (2.5 \pm 4.3 \pm 0.7) \times 10^{-4}$.

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I. INTRODUCTION

Charge-parity (CP) violation is one of the key ingredients that are needed to generate the asymmetry between matter and antimatter observed in the Universe [1]. The Standard Model (SM) of particle physics, where all known CP -violating processes arise from the irreducible phase of the Cabibbo-Kobayashi-Maskawa matrix [2,3], is, however, unable to explain the observed asymmetry [4,5]. New dynamics that lead to a significant enhancement of CP -violating processes are required, making searches for CP violation a powerful probe for physics beyond the SM. Although CP violation has been experimentally observed in the down-type quark sector with measurements of K and B mesons [6–10], no indication of new dynamics has been reported yet. Only recently has CP violation been observed in the decay of charmed mesons [11]. The limited precision of the SM predictions, together with the limited amount of experimental information available [12], is, however, not yet sufficient to establish whether the observed signal could be explained by the SM [13–18]. Additional searches for CP violation in the charm sector, and particularly for more suppressed and yet-to-be-observed signs of CP -violating effects induced by D^0 - \bar{D}^0 mixing, have unique potential to probe for the existence of beyond-the-SM dynamics, which couple preferentially to up-type quarks [19–24].

This paper reports a search for CP violation in D^0 - \bar{D}^0 mixing, or in the interference between mixing and decay, through the measurement of the asymmetry between the effective decay widths, $\hat{\Gamma}$, of mesons initially produced as D^0 and \bar{D}^0 and decaying into the CP -even final states $f = K^+ K^-$, $\pi^+ \pi^-$:

$$A_\Gamma(f) \equiv \frac{\hat{\Gamma}(D^0 \rightarrow f) - \hat{\Gamma}(\bar{D}^0 \rightarrow f)}{\hat{\Gamma}(D^0 \rightarrow f) + \hat{\Gamma}(\bar{D}^0 \rightarrow f)}. \quad (1)$$

Several measurements of the parameter $A_\Gamma(f)$ have been performed by the BABAR [25], CDF [26], Belle [27], and LHCb [28–30] Collaborations, leading to the current world-average value of $(-3.2 \pm 2.6) \times 10^{-4}$ [12], when neglecting differences between the $D^0 \rightarrow K^+ K^-$ and $D^0 \rightarrow \pi^+ \pi^-$ decays.¹ The achieved sensitivity is still 1 order of magnitude larger than the theoretical predictions of $A_\Gamma \approx 3 \times 10^{-5}$ [31]. This paper updates the LHCb measurements of Refs. [28–30] using the data sample of proton-proton collisions collected at a center-of-mass energy of 13 TeV during 2016–2018, and corresponding to an integrated luminosity of 5.4 fb^{-1} . The analysis is performed using D^0 mesons originating from semileptonic decays of b hadrons, where the b -hadron candidates are only partially reconstructed. The charge of the muon identifies (“tags”) the flavor of the D^0 meson at its production. The samples are dominated by $B^- \rightarrow D^0 \mu^- X$ and $\bar{B}^0 \rightarrow D^0 \mu^- X$ decays,

¹Throughout the paper, the inclusion of the charge-conjugate decay mode is implied unless otherwise stated.

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where X denotes any set of final-state particles that are not reconstructed.

The paper is structured as follows: the analysis strategy is described in Sec. II; the LHCb detector is sketched in Sec. III; Sec. IV details the criteria used to select the signal and control samples; Sec. V describes the fit method, and its validation using $D^0 \rightarrow K^- \pi^+$ decays; the determination of the systematic uncertainties is outlined in Sec. VI, before concluding with the presentation of the final results in Sec. VII.

II. ANALYSIS STRATEGY

Due to the weak interactions, the mass eigenstates of neutral charm mesons, D_1 and D_2 , are a superposition of the flavor states, D^0 and \bar{D}^0 : $|D_{1,2}\rangle \equiv p|D^0\rangle \pm q|\bar{D}^0\rangle$, where q and p are complex coefficients satisfying $|p|^2 + |q|^2 = 1$. Hence, an originally produced D^0 meson can oscillate as a function of time into a \bar{D}^0 meson, and vice versa, before decaying. In the limit of CP symmetry, q equals p and the oscillations are characterized by only two dimensionless parameters, $x \equiv (m_1 - m_2)c^2/\Gamma$ and $y \equiv (\Gamma_1 - \Gamma_2)/2\Gamma$, where $m_{1(2)}$ and $\Gamma_{1(2)}$ are the mass and decay width of the CP -even (odd) eigenstate $D_{1(2)}$, respectively, and $\Gamma \equiv (\Gamma_1 + \Gamma_2)/2$ is the average decay width [32]. The values of x and y have been measured to be of the order of 1% or smaller [12]. In the presence of CP violation, the mixing rates for mesons produced as D^0 and \bar{D}^0 differ, further enriching the phenomenology. As an example, indicating with A_f (\bar{A}_f) the decay amplitude of a D^0 (\bar{D}^0) meson into the final state f , three different manifestations of CP violation can be measured: (i) CP violation in the decay if $\mathcal{A}_{CP}^{\text{dir}}(f) \equiv (|A_f|^2 - |\bar{A}_f|^2)/(|A_f|^2 + |\bar{A}_f|^2)$ differs from zero, (ii) CP violation in mixing if $|q/p|$ differs from unity, and (iii) CP violation in the interference between mixing and decay if $\phi_f \equiv \arg[(q\bar{A}_f)/(pA_f)]$ differs from zero. The latter two can be accessed by measuring the decay-time-dependent CP asymmetry

$$\mathcal{A}_{CP}(D^0 \rightarrow f; t) = \frac{\Gamma(D^0(t) \rightarrow f) - \Gamma(\bar{D}^0(t) \rightarrow f)}{\Gamma(D^0(t) \rightarrow f) + \Gamma(\bar{D}^0(t) \rightarrow f)}. \quad (2)$$

In the limit of small mixing parameters, Eq. (2) can be approximated as a linear function of decay time [33,34],

$$\mathcal{A}_{CP}(D^0 \rightarrow f; t) \approx \mathcal{A}_{CP}^{\text{dir}}(f) - A_\Gamma(f) \frac{t}{\tau}, \quad (3)$$

where $\tau = 1/\Gamma$ is the average lifetime of neutral D mesons. The coefficient $A_\Gamma(f)$ is related to the mixing and CP -violation parameters by [35]

$$A_\Gamma(f) \approx -x\phi_f + y(|q/p| - 1) - y\mathcal{A}_{CP}^{\text{dir}}(f). \quad (4)$$

Contrarily to the measurement reported in Ref. [11], which is sensitive to $\mathcal{A}_{CP}^{\text{dir}}(K^+K^-) - \mathcal{A}_{CP}^{\text{dir}}(\pi^+\pi^-)$, $A_\Gamma(f)$ is mostly

sensitive to CP violation in mixing or in the interference between mixing and decay, because the term $y\mathcal{A}_{CP}^{\text{dir}}(f) \leq 10^{-5}$ [12] can be neglected at the current level of experimental precision. Moreover, neglecting the $\mathcal{O}(10^{-3})$ difference between the weak phases of the decay amplitudes to the CP -even final states K^+K^- and $\pi^+\pi^-$, $\phi_f \approx \phi \equiv \arg(q/p)$ becomes universal and A_Γ independent of f [22].

Experimentally, the partial rate asymmetry of Eq. (2) cannot be measured directly because of charge-asymmetric detection efficiencies and asymmetric production rates of D^0 and \bar{D}^0 mesons from semileptonic b -hadron decays in proton-proton collisions. Instead, the “raw” asymmetry between the D^0 and \bar{D}^0 mesons yields,

$$A_{\text{raw}}(D^0 \rightarrow f) = \frac{N(\bar{B} \rightarrow D^0(\rightarrow f)\mu^-X) - N(B \rightarrow \bar{D}^0(\rightarrow f)\mu^+X)}{N(\bar{B} \rightarrow D^0(\rightarrow f)\mu^-X) + N(B \rightarrow \bar{D}^0(\rightarrow f)\mu^+X)}, \quad (5)$$

is measured as a function of decay time. Neglecting higher-order terms in the involved asymmetries, which are at most $\mathcal{O}(1\%)$, the raw asymmetry can be approximated as

$$A_{\text{raw}}(D^0 \rightarrow f; t) \approx \mathcal{A}_{CP}(D^0 \rightarrow f; t) + A_D(\mu) + A_P(D), \quad (6)$$

where $A_D(\mu)$ and $A_P(D)$ are the nuisance asymmetries due to the detection efficiency of the tagging muon and to the production rates of the neutral D mesons, respectively. The parameter A_Γ corresponds to the slope of the decay-time-dependent raw asymmetry only if A_D and A_P are independent of decay time. In this analysis, a possible time dependence of A_D and A_P is considered as a source of systematic uncertainty. The analysis procedure is validated on data using a control sample of Cabibbo-favored $D^0 \rightarrow K^- \pi^+$ decays, whose size exceeds that of the $D^0 \rightarrow K^+K^-$ and $D^0 \rightarrow \pi^+\pi^-$ signal modes by approximately 1 order of magnitude, and where measured asymmetries can be attributed solely to instrumental effects because no CP violation is expected. To avoid potential experimenter's bias, the measured values of $A_\Gamma(K^+K^-)$ and $A_\Gamma(\pi^+\pi^-)$ remained unknown during the development of the analysis and were examined only after the analysis procedure and the evaluation of the systematic uncertainties were finalized.

III. DETECTOR

The LHCb detector [36,37] is a single-arm forward spectrometer covering the pseudorapidity range $2 < \eta < 5$, designed for the study of particles containing b or c quarks. The detector includes a high-precision tracking system consisting of a silicon-strip vertex detector surrounding the pp interaction region, a large-area silicon-strip detector located upstream of a dipole magnet with a bending power of about 4 Tm, and three stations of silicon-strip detectors and straw drift tubes placed downstream of the magnet.

The tracking system provides a measurement of the momentum, p , of charged particles with relative uncertainty that varies from 0.5% at low momentum to 1.0% at 200 GeV/ c . The minimum distance of a track to a primary vertex (PV), the impact parameter, is measured with a resolution of $(15 + 29/p_T) \mu\text{m}$, where p_T is the component of the momentum transverse to the beam, in GeV/ c . Different types of charged hadrons are distinguished using information from two ring-imaging Cherenkov detectors. Photons, electrons, and hadrons are identified by a calorimeter system consisting of scintillating-pad and preshower detectors, an electromagnetic and a hadronic calorimeter. Muons are identified by a system composed of alternating layers of iron and multiwire proportional chambers. The magnetic-field polarity is reversed periodically during data taking to mitigate the differences of reconstruction efficiencies of particles with opposite charges.

The on-line event selection is performed by a trigger, which consists of a hardware stage followed by a two-level software stage. In between the two software stages, an alignment and calibration of the detector is performed in near real time [38]. The same alignment and calibration information is propagated to the off-line reconstruction, ensuring consistent and high-quality particle identification information between the trigger and off-line software. The identical performance of the on-line and off-line reconstructions offers the opportunity to perform physics analyses directly using candidates reconstructed in the trigger [39,40], which the present analysis exploits.

IV. SELECTION

The selection criteria are mainly inherited from the measurement of the difference between the decay-time-integrated CP asymmetries in $D^0 \rightarrow K^+K^-$ and $D^0 \rightarrow \pi^+\pi^-$ decays [11], which uses the same sample of proton-proton collisions. Signal candidates are first required to pass the hardware trigger, which selects events containing at least one charged particle with high transverse momentum that leaves a track in the muon system. At the first stage of the software trigger, events are selected if they contain at least one track having large transverse momentum and being incompatible with originating from any PV, or if any two-track combination forming a secondary vertex passes a multivariate classifier. If a particle is identified as a muon, a lower p_T threshold is applied. At the second stage of the software trigger, the full event reconstruction is performed, and requirements on kinematic, topological, and particle-identification criteria are placed on the signal candidates. A D^0 candidate is formed by combining two well-reconstructed, oppositely charged tracks such that they are consistent with originating from a common vertex. The D^0 candidate must satisfy requirements on the vertex quality and has to be well separated from all PVs in the event. At the next step, the D^0 candidate is combined with a muon to form a B candidate. Only candidates where the D^0

meson decays downstream along the beam axis with respect to the B candidate are further considered. The B candidate must have a visible mass, $m(D^0\mu)$, and a corrected mass, $m_{\text{corr}}(B)$, consistent with a signal decay. The corrected mass is computed as $m_{\text{corr}}(B) \equiv \sqrt{m^2(D^0\mu) + p_\perp^2(D^0\mu)} + p_\perp(D^0\mu)$, where $p_\perp(D^0\mu)$ is the momentum of the $D^0\mu$ system transverse to the B flight direction, to partially correct for the unreconstructed particles in the decay of the B hadron.

In the off-line selection, trigger signals are associated with reconstructed particles. Particle-identification criteria and requirements on $m(D^0\mu)$ and $m_{\text{corr}}(B)$ are tightened with respect to the on-line selection. The mass of the D^0 candidate is required to be in the ranges [1825, 1925] MeV/ c^2 , [1820, 1939] MeV/ c^2 , and [1780, 1940] MeV/ c^2 for $D^0 \rightarrow K^+K^-$, $D^0 \rightarrow \pi^+\pi^-$, and $D^0 \rightarrow K^-\pi^+$ decays, respectively, to reduce the amount of background decays with misidentified final-state particles to a negligible level. The reconstructed decay time is computed from the distance, L , between the measured D^0 and B decay vertices and from the D^0 momentum, $p(D^0)$, as $t = m_{D^0}L/[p(D^0)c]$, where m_{D^0} is the known D^0 mass [32]. All D^0 candidates with a reconstructed decay time that is either negative or exceeds 10 times the D^0 lifetime are discarded. Mass vetoes suppress background from misreconstructed B decays to final states involving a charmonium resonance, such as $B^- \rightarrow \psi^{(\prime)}(\rightarrow \mu^+\mu^-)h^-$ with $h = \pi$ or K , where a muon is misidentified as a pion or kaon and is used in the D^0 final state. Tag muons reconstructed in regions of phase space with large instrumental asymmetries, due to muons of one charge either being bent out of the detector acceptance or deflected into the LHC beam pipe, are vetoed. The fraction of signal candidates removed by this requirement is 10%. In addition, for $D^0 \rightarrow K^-\pi^+$ decays, candidates with kaon $p_T < 800$ MeV/ c are removed to reduce instrumental asymmetry between the detection of negatively and positively charged kaons. Since these requirements do not reduce the background to a sufficiently low level for $D^0 \rightarrow K^+K^-$ and $D^0 \rightarrow \pi^+\pi^-$ decays, a dedicated boosted decision tree (BDT) is trained to isolate the signal candidates from background made of accidental combinations of charged particles (“combinatorial background”). The variables used in the BDT to discriminate signal from combinatorial background are the fit quality of the D^0 and the B decay vertices, the D^0 flight distance, the D^0 impact parameter with respect to the closest PV, the transverse momenta of the D^0 decay products, the significance of the distance between the D^0 and B decay vertices, and the visible and corrected masses of the B -hadron candidate. The BDT is trained using $D^0 \rightarrow K^-\pi^+$ decays as signal proxies and candidates from the D^0 mass sidebands of the signal decay modes as background. The optimal requirement on the BDT discriminant is chosen by maximizing the figure of

merit $\mathcal{S}/\sqrt{\mathcal{S}+\mathcal{B}}$ in a range corresponding to approximately 3 times the mass resolution around the D^0 mass, where \mathcal{S} and \mathcal{B} denote the signal and background yields, respectively. If an event contains more than one candidate after the full selection, one is chosen at random. The fraction of candidates removed by this requirement is 0.4%.

The mass distributions of the selected signal- and control-decay candidates are shown in Fig. 1. Details about the fit model are given in the next section. Approximately 9×10^6 , 3×10^6 , and 76×10^6 signal $D^0 \rightarrow K^+K^-$, $D^0 \rightarrow \pi^+\pi^-$, and $D^0 \rightarrow K^-\pi^+$ decays, respectively, are

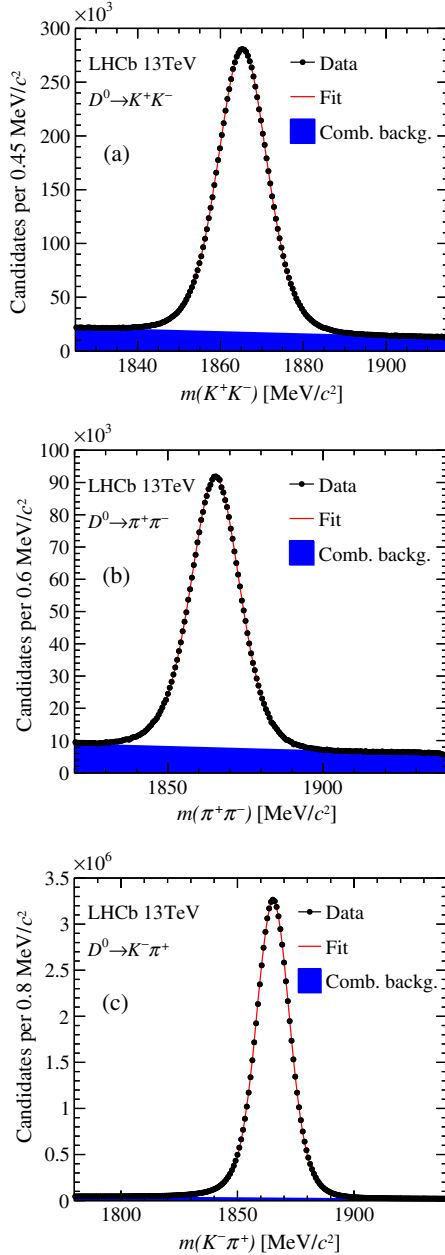


FIG. 1. Mass distributions of (a) $D^0 \rightarrow K^+K^-$, (b) $D^0 \rightarrow \pi^+\pi^-$, and (c) $D^0 \rightarrow K^-\pi^+$ candidates with fit projections overlaid.

reconstructed over a smooth background dominated by accidental combinations of charged particles.

V. FIT METHOD

The samples of selected $D^0 \rightarrow K^+K^-$, $D^0 \rightarrow \pi^+\pi^-$, and $D^0 \rightarrow K^-\pi^+$ candidates are split into 20 approximately equally populated subsets (“bins”) of decay time in the range $[0, 10]\tau$. In each decay-time bin, the raw asymmetry A_{raw} is determined by a simultaneous binned χ^2 fit to the $m(D^0)$ distributions of the D^0 and \bar{D}^0 candidates, split according to the muon tag. The total signal yields and asymmetries are treated as shared floating parameters of the fit. The fits include two components: signal and combinatorial background. The signal is described with a sum of a Gaussian and a Johnson’s S_U distribution [41], with parameters determined from a fit to the decay-time-integrated mass spectra. To account for the observed dependence of the signal mass shape on decay time, the means and widths of the signal distributions are left free to float individually for each decay-time bin. The mass shape is assumed to be the same for D^0 and \bar{D}^0 candidates for charge-symmetric final states of the signal modes, and allowed to differ for $D^0 \rightarrow K^-\pi^+$ and $\bar{D}^0 \rightarrow K^+\pi^-$ candidates. The combinatorial background is described by a linear function, with a slope that floats independently in each decay-time bin and is allowed to differ between D^0 and \bar{D}^0 candidates.

The raw asymmetry measured in decay-time bin i is fit by minimizing the least squares with respect to the linear function $A_{\text{raw}}(0) - A_{\Gamma}\langle t \rangle_i/\tau$. The decay-time-independent terms of Eqs. (3) and (6) are incorporated into a single parameter, $A_{\text{raw}}(0)$, that is determined by the fit together with A_{Γ} . The average decay time in each bin i , $\langle t \rangle_i$, is computed using the decay-time distribution of background-subtracted D^0 candidates. Statistically consistent values are found for the control and signal modes. The D^0 lifetime τ is set to its known value [32]. Using large samples of simulated experiments, it is verified that the analysis procedure leads to unbiased estimates of the fit parameters and of their uncertainties. Figure 2 shows the projection of the decay-time-dependent fit to the $D^0 \rightarrow K^-\pi^+$ control

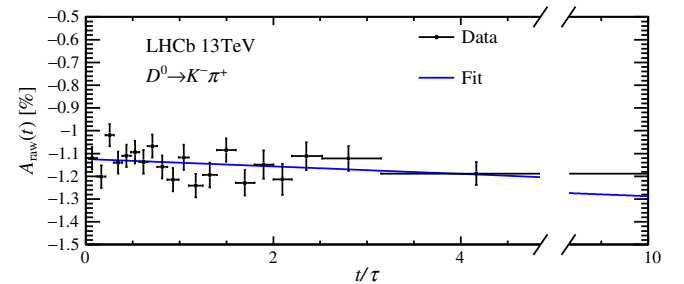


FIG. 2. Raw asymmetry as a function of decay time with fit projection overlaid for $D^0 \rightarrow K^-\pi^+$ signal candidates.

sample. Here A_Γ is measured to be $(1.6 \pm 1.2) \times 10^{-4}$, where the uncertainty is statistical only. The measured value is consistent with zero as expected, confirming the validity of the assumption of decay-time-independent nuisance asymmetries. In $D^0 \rightarrow K^- \pi^+$ decays, due to their charge-asymmetric final states, detection asymmetries are more pronounced compared to the signal modes, where these asymmetries are only caused by the muons used to tag the flavor of the D^0 mesons.

VI. SYSTEMATIC UNCERTAINTIES

The systematic uncertainty is dominated by the following contributions: the impact of decay-time acceptance and resolution, the effect of neglected background from combinations of real D^0 candidates with unrelated muons (which might lead to a wrong identification of the neutral D -meson flavor), and the impact of the assumed parametrization of the signal and background mass shapes. These effects are studied using large samples of pseudoexperiments, where the above sources of systematic biases are simulated.

The average decay-time resolution is estimated to be 127 fs using simulated decays. In the generation of the pseudoexperiments, the resolution is increased by 10% to account for differences between data and simulation. The decay-time acceptance is estimated from data by comparing the background-subtracted decay-time distributions of $D^0 \rightarrow K^- \pi^+$ candidates with an exponential function convoluted with the decay-time resolution. Different sets of pseudoexperiments, simulating the effect of decay-time acceptance and resolution, are generated with values of A_Γ in the range $[-30, 30] \times 10^{-4}$. Each pseudoexperiment is then fit with the default analysis approach, and the difference between the measured and the input values of A_Γ is used to determine the systematic bias. As the bias is found to depend linearly on the true value of A_Γ , the largest bias observed within the 68% confidence-level interval of the current world average [12] is taken as the systematic uncertainty. This amounts to 0.3×10^{-4} (0.4×10^{-4}) for $D^0 \rightarrow K^+ K^-$ ($D^0 \rightarrow \pi^+ \pi^-$) decays.

The probability to wrongly associate unrelated muons with the D^0 candidates is estimated using the yields of “wrong-sign” $D^0(\rightarrow K^- \pi^+) \mu^+$ and $\bar{D}^0(\rightarrow K^+ \pi^-) \mu^-$ candidates in data, which are corrected for the rate of doubly Cabibbo-suppressed decays and decays due to flavor oscillation using the measurements reported in Ref. [42]. Mistag probabilities between 1% at low decay times and 3% at high decay times are observed. Also, in this case, the bias observed in pseudoexperiments depends linearly on the true value of A_Γ . Following the same strategy as discussed above, a systematic uncertainty of 0.3×10^{-4} (0.6×10^{-4}) is assigned for $D^0 \rightarrow K^+ K^-$ ($D^0 \rightarrow \pi^+ \pi^-$) decays.

To estimate any potential bias due to the specific choice of the mass model used in the fits that determine the raw

TABLE I. Summary of the dominant contributions to the systematic uncertainty on $A_\Gamma(K^+ K^-)$ and $A_\Gamma(\pi^+ \pi^-)$.

Source of uncertainty	$A_\Gamma(K^+ K^-)$ [10^{-4}]	$A_\Gamma(\pi^+ \pi^-)$ [10^{-4}]
Decay-time resolution and acceptance	0.3	0.4
Mistag probability	0.3	0.6
Mass-fit model	0.3	0.3
Total	0.5	0.8

asymmetries, samples of pseudoexperiments are generated using alternative signal and background models that describe the data equally well. The observed bias is independent of the input A_Γ and results in an additional systematic uncertainty of 0.3×10^{-4} for both signal decay channels.

Uncertainties on $\langle t \rangle_i / \tau$ arising from relative misalignments of subdetectors and from the uncertainty on the input value of the D^0 lifetime [32] give negligible contributions. Furthermore, unexpected biases due to a possible decay-time dependence of the nuisance asymmetries and due to the selection procedure are investigated using the $D^0 \rightarrow K^- \pi^+$ control sample and/or by measuring A_Γ in disjoint subsamples split by magnetic-field polarity, year of data taking, and kinematic variables of the B hadron, D^0 meson, and muon candidates. No unexpected variations are observed, and no additional systematic uncertainties are assigned.

A summary of the relevant systematic uncertainties is given in Table I. The total systematic uncertainty is obtained by summing in quadrature the individual components and amounts to 0.5×10^{-4} and 0.8×10^{-4} for $A_\Gamma(K^+ K^-)$ and $A_\Gamma(\pi^+ \pi^-)$, respectively.

VII. RESULTS AND CONCLUSIONS

A search for decay-time-dependent CP violation in $D^0 \rightarrow K^+ K^-$ and $D^0 \rightarrow \pi^+ \pi^-$ decays is performed using proton-proton collision data recorded with the LHCb detector at a center-of-mass energy of 13 TeV, and corresponding to an integrated luminosity of 5.4 fb^{-1} . The D^0 mesons are required to originate from semileptonic b -hadron decays, such that the charge of the muon identifies the flavor of the neutral D meson at the moment of its production. The parameter A_Γ is determined from a fit to the asymmetry between D^0 and \bar{D}^0 yields as a function of decay time. The projections of the fits for both $D^0 \rightarrow K^+ K^-$ and $D^0 \rightarrow \pi^+ \pi^-$ samples are shown in Fig. 3. The results are

$$A_\Gamma(K^+ K^-) = (-4.3 \pm 3.6 \pm 0.5) \times 10^{-4},$$

$$A_\Gamma(\pi^+ \pi^-) = (2.2 \pm 7.0 \pm 0.8) \times 10^{-4},$$

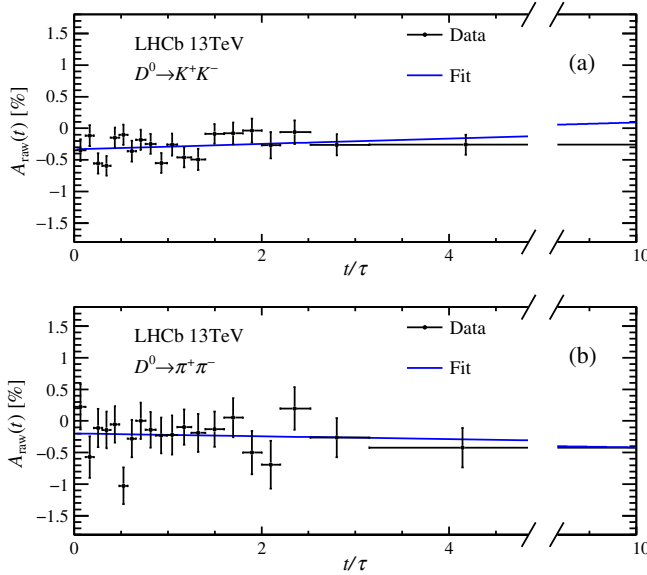


FIG. 3. Raw asymmetry as a function of decay time with fit projection overlaid for (a) $D^0 \rightarrow K^+ K^-$ and (b) $D^0 \rightarrow \pi^+ \pi^-$ signal candidates.

where the uncertainties are statistical and systematic, respectively.

The measured values are combined with previous LHCb measurements based on data corresponding to 3 fb^{-1} collected at center-of-mass energies of 7 and 8 TeV, and where the neutral D mesons originate either from semi-leptonic b -hadron decays [28] or from promptly produced $D^{*+}(2010)$ mesons [29], with which they are consistent. The combination accounts for correlations in the systematic uncertainties and yields

$$A_{\Gamma}(K^+ K^-) = (-4.4 \pm 2.3 \pm 0.6) \times 10^{-4},$$

$$A_{\Gamma}(\pi^+ \pi^-) = (2.5 \pm 4.3 \pm 0.7) \times 10^{-4}.$$

Assuming A_{Γ} to be universal, the above two results can be averaged to yield $A_{\Gamma} = (-2.9 \pm 2.0 \pm 0.6) \times 10^{-4}$. The results do not show any indication of CP violation in charm mixing or in the interference between mixing and decay.

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